

1a)

Writing differential form of potential  $dV = \frac{dq}{4\pi\epsilon_0 r}$ . (2 point)

Drawing a diagram, labelling distance to each point  $r = \sqrt{z^2 + R^2}$ , constant. (2 point)

Since we are given total charge  $Q$  and not linear charge density, we can do the integral simply

$$V = \frac{1}{4\pi\epsilon_0} \frac{1}{\sqrt{z^2 + R^2}} \int dq = \frac{Q}{4\pi\epsilon_0} \frac{1}{\sqrt{z^2 + R^2}}. \quad (4 \text{ point})$$

1b)

Electric field given by  $\vec{E} = E_0 \frac{\exp(-\kappa r)}{r^2} \hat{r}$ . Note that dimensions of  $[E_0] = [q/\epsilon_0]$ . (1 point)

To find charge enclosed in radius  $1/\kappa$ , use Gauss's law and symmetry. (3 point)

Choosing gaussian sphere with radius  $1/\kappa$  (2 point), we can immediately solve for  $Q_{enc}$

$$(E_0 e^{-1} \kappa^2) \frac{4\pi}{\kappa^2} = \frac{Q_{enc}}{\epsilon_0} \Rightarrow Q_{enc} = E_0 (4\pi\epsilon_0) e^{-1}. \quad (2 \text{ point})$$

1c)

Multiple ways to do this. All will get +2 points.

Use differential form of Gauss's law  $\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$ .

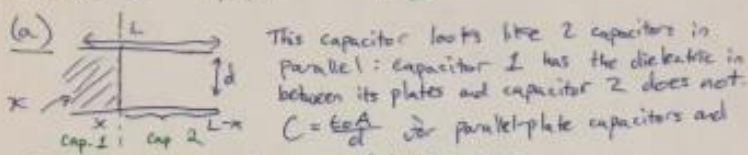
Plugging in for the given  $\vec{E} = E(r)\hat{r}$ , we find that (in spherical coordinates),

$$\frac{\rho}{\epsilon_0} = \vec{\nabla} \cdot \vec{E} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 E(r)) = \frac{2E(r)}{r} + E'(r) = -E_0 \kappa \frac{\exp(-\kappa r)}{r^2}$$

PHYS 7B, Section 3 (Wurtele)

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Midterm #2 4/3/17 #2



$C = \frac{\epsilon_0 A}{d}$  for parallel-plate capacitors and  $C = \kappa C_0$  for capacitors with dielectrics, so:

$$C_1 = \kappa \frac{\epsilon_0 x}{d}, \quad C_2 = \frac{\epsilon_0 (L-x)}{d}$$

Since capacitors add in parallel, the total capacitance is

$$C = C_1 + C_2 = \frac{\epsilon_0 L}{d} [\kappa x + L - x]$$

$$(b) U = \frac{1}{2} \frac{Q^2}{C} = \frac{1}{2} \frac{(CV)^2}{C} = \frac{1}{2} CV^2$$

$$\Rightarrow U = \frac{1}{2} \frac{\epsilon_0 L}{d} [\kappa x + L - x] V^2$$

(c) After the voltage is disconnected, the circuit is broken so that charges cannot flow anymore. So charge does not change if the dielectric is removed.

If the dielectric is removed, the capacitance is just  $\frac{\epsilon_0 L}{d}$  (it looks like a single capacitor). Thus, the final voltage is

$$V_f = \frac{Qd}{\epsilon_0 L} \quad \text{Also, } Q = \frac{\epsilon_0 L}{d} [\kappa x + L - x] V$$

$$\Rightarrow V_f = V \left[ \frac{\kappa x + L - x}{L} \right] = V \left[ 1 + \frac{\kappa}{L} (x - L) \right]$$

$\Delta V = V_f - V = V \frac{\kappa}{L} (x - L) > 0$  (since  $\kappa > 1$ ), so voltage increases.

$$U_f = \frac{1}{2} C_f V_f^2 = \frac{1}{2} \frac{\epsilon_0 L}{d} V^2 \left[ \frac{\kappa x + L - x}{L} \right]^2 = \frac{1}{2} \frac{\epsilon_0 L}{d} [\kappa x + L - x] V^2 \left[ 1 + \frac{\kappa}{L} (x - L) \right]$$

$\Delta U = U_f - U_i = \frac{1}{2} \frac{\epsilon_0 \kappa}{d} V^2 (x - L) > 0$ , so energy increases.

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(d) Since the voltage source is connected, the voltage does not change

if the dielectric is removed. Again,  $C_f = \frac{\epsilon_0 L}{d}$  if the dielectric is removed

$$\text{so: } Q_f = C_f V = \frac{\epsilon_0 L}{d} V, \quad Q_i = \frac{\epsilon_0 L}{d} [\kappa x + L - x] V$$

$$= Q_i \cdot \frac{L}{L + \kappa x - x}$$

$\Delta Q = Q_f - Q_i = \frac{\epsilon_0 L}{d} x [1 - \kappa] V < 0$  (since  $\kappa > 1$ ), so charge decreases

$$U_f = \frac{1}{2} C_f V^2 = \frac{1}{2} \frac{\epsilon_0 L}{d} V^2 = U_i \left[ \frac{L}{L + \kappa x - x} \right]$$

$$U_i = \frac{1}{2} \frac{\epsilon_0 L}{d} [\kappa x + L - x] V^2$$

$\Delta U = U_f - U_i = \frac{1}{2} \frac{\epsilon_0 L}{d} x [1 - \kappa] V^2 < 0$  (since  $\kappa > 1$ ), so energy decreases.

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**Problem 3a.** Find the electric field at a point  $x$  along the line joining the wires.

**Solution** The electric field outside a cylinder of uniform charge  $Q$  and length  $l$  is:

$$\mathbf{E} = \frac{Q}{2\pi\epsilon_0 l r} \hat{r} \quad (1)$$

And thus at a point  $x$  between the wires, the net electric field is:

$$\mathbf{E}_{net} = \frac{-Q}{2\pi\epsilon_0 l} \left( \frac{1}{x} + \frac{1}{d-x} \right) \hat{x} \quad (2)$$

Note both fields point in the negative  $\hat{x}$  (towards the negatively charged wire) at  $x$ .

**Problem 3b.** Find the potential at  $x$ .

**Solution** In general,

$$V(a) - V(b) = - \int_b^a \mathbf{E} \cdot d\mathbf{s}$$

We can't choose  $V(\infty)$  to be finite since it diverges at infinity. Technically one would need to use a different electric field equation than (3) to show this since (3) is only valid for  $R < x < d - R$ . Regardless, the real electric field line integral diverges at infinity, so our problem remains. Thus, let us choose  $b = R$  and  $a = x$ :

$$V(x) = V(R) - \int_R^x \frac{-Q}{2\pi\epsilon_0 l} \left( \frac{1}{x'} + \frac{1}{d-x'} \right) dx' \quad (3)$$

$$= V(R) + \frac{Q}{2\pi\epsilon_0 l} \ln \left( \frac{x(d-R)}{R(d-x)} \right) \quad (4)$$

**Problem 3c.** Find the capacitance per unit length of the two wires.

**Solution** The capacitance is

$$C = \frac{Q}{V} \quad (5)$$

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Where  $V$  is the voltage difference between the two wires, that is  $V(d-R) - V(R)$ . Using our result from part *b*, we have:

$$V \equiv V(d-R) - V(R) \quad (6)$$

$$= \frac{Q}{2\pi\epsilon_0 l} \ln \left( \frac{(d-R)^2}{R^2} \right) \quad (7)$$

$$= \frac{Q}{\pi\epsilon_0 l} \ln \left( \frac{d-R}{R} \right) \quad (8)$$

And consequently, the capacitance is:

$$C = \frac{Q}{V} \quad (9)$$

$$= \frac{Q}{\frac{Q}{\pi\epsilon_0 l} \ln \left( \frac{d-R}{R} \right)} \quad (10)$$

$$= \frac{\pi\epsilon_0 l}{\ln \left( \frac{d-R}{R} \right)} \quad (11)$$

Just divide both sides by  $l$  to obtain the capacitance per unit length.

Problem 4

a)  $R = r_0 (1 + \frac{z}{L})$  think of each cylindrical slice as a resistor in series  $\rightarrow$  their resistances add linearly

$$R = \int dR = \int_0^L \frac{\rho dz}{A} = \int_0^L \frac{\rho dz}{\pi r_0^2 (1 + \frac{z}{L})^2} = \frac{\rho L}{\pi r_0^2} \left( \frac{1}{L} - \frac{1}{2L} \right) = \frac{\rho L}{2\pi r_0^2}$$

b)  $P = I^2 R = I^2 \frac{\rho L}{2\pi r_0^2}$

c) Looking at a single cylindrical slice of length  $dz$  its power dissipation is  $\frac{I^2 \rho dz}{\pi r_0^2 (1 + \frac{z}{L})^2}$  so  $\frac{dP}{dz} = \frac{I^2 \rho}{\pi r_0^2 (1 + \frac{z}{L})^2}$

d) Consider a volume of area  $dA$  and length  $dz$  through which a current density  $J$  passes through

The power dissipated in this volume is  $(J dA)^2 \frac{\rho dz}{dA} = J^2 \rho dV$

$$\text{So } \frac{dP}{dV} = J^2 \rho$$


rubric

a) mention the slices are in series  $\rightarrow 4$

Set up  $R$  integral correctly  $\rightarrow 4$

Evaluate the integral correctly  $\rightarrow 4$


b, c, d answer correct full points (won't double penalize mistakes)

6) a) forces cancel pairwise (9 o'clock balances 7 o'clock)  
 $F = 0$  

b)  $W = qV = \frac{12 q Q}{4\pi \epsilon_0 R} = \frac{3 q Q}{\epsilon_0 R}$

c) all forces cancel pairwise except '1 o'clock' charge  
 $L \theta = 0 \quad \theta = 2\pi$

c) all forces cancel pairwise except '1 o'clock' charge

  $\theta = \frac{2\pi}{12}$   $F = \frac{qQ}{4\pi \epsilon_0 R^2}$

