

Superposition:  $V(\vec{r}) = \int \frac{\rho(\vec{r}') d\vec{r}'}{4\pi\epsilon_0 |\vec{r} - \vec{r}'|}$

$\vec{r} = (x, 0, 0)$ ,  $\vec{r}' = (0, a\cos\theta, a\sin\theta)$

$\vec{r} - \vec{r}' = (x, -a\cos\theta, -a\sin\theta)$

$|\vec{r} - \vec{r}'| = \sqrt{x^2 + a^2}$

Potential from ring

$\Rightarrow V(x, 0, 0) = \int \frac{\rho(\vec{r}') d\vec{r}'}{4\pi\epsilon_0 \sqrt{x^2 + a^2}} \rightarrow \int \frac{\lambda r d\theta}{4\pi\epsilon_0 \sqrt{x^2 + a^2}} = \frac{\lambda a z \pi}{4\pi\epsilon_0 \sqrt{x^2 + a^2}} = \frac{\lambda a}{2\epsilon_0 \sqrt{x^2 + a^2}}$

$\lambda = \frac{q}{2\pi a}$

Potential from point charge

$V_p(x, 0, 0) = \frac{-q}{4\pi\epsilon_0 x}$

$\Rightarrow V_{tot}(x, 0, 0) = \frac{q}{4\pi\epsilon_0} \left[ \frac{-1}{x} + \frac{1}{\sqrt{x^2 + a^2}} \right]$

b)  $\vec{E} = -\vec{\nabla}\phi$ .  $\phi$  depends only on  $x$ , so  $\vec{E}$  has only an  $x$ -direction.

$\vec{E} = \frac{-q}{4\pi\epsilon_0} \left[ \frac{\partial}{\partial x} (x^{-1}) + \frac{\partial}{\partial x} (x^2 + a^2)^{-1/2} \right] \hat{x} = \frac{-q}{4\pi\epsilon_0} \left[ \frac{+1}{x^2} - \frac{1(2x)}{2(x^2 + a^2)^{3/2}} \right] \hat{x}$

$\vec{E}(x, 0, 0) = \frac{q}{4\pi\epsilon_0} \left[ \frac{-1}{x^2} + \frac{x}{(x^2 + a^2)^{3/2}} \right] \hat{x}$

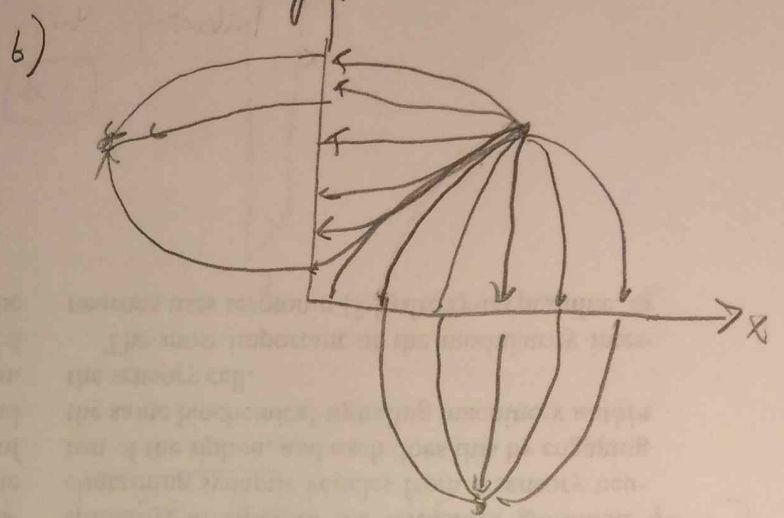
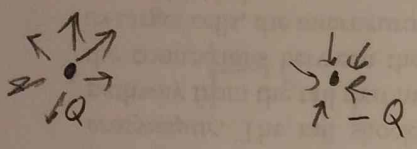
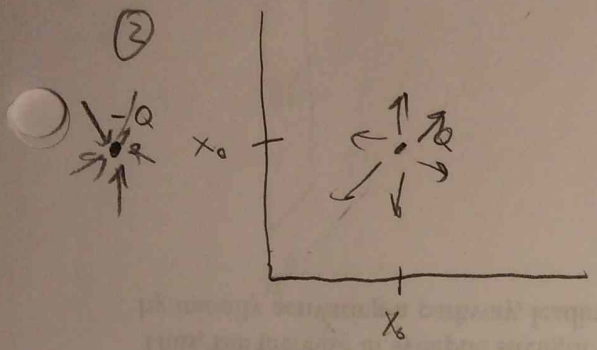
(c) Want to take  $x \gg a$ . This always suggests the use of a Taylor Expansion.

$\frac{x}{(x^2 + a^2)^{3/2}} = \frac{x}{x^3 (1 + \frac{a^2}{x^2})^{3/2}} = \frac{1}{x^2} \left( 1 - \frac{3}{2} \left( \frac{a}{x} \right)^2 + o\left( \left( \frac{a}{x} \right)^4 \right) \right)$

$\vec{E}(x, 0, 0) = \frac{q}{4\pi\epsilon_0} \hat{x} \left[ \frac{-1}{x^2} + \frac{1}{x^2} - \frac{3}{2x^2} \left( \frac{a}{x} \right)^2 + o\left( \left( \frac{a}{x} \right)^4 \right) \right]$

$\vec{E}(x, 0, 0) = \frac{q}{4\pi\epsilon_0} \hat{x} \left[ -\frac{3}{2x^2} \left( \frac{a}{x} \right)^2 + o\left( \left( \frac{a}{x} \right)^4 \right) \right]$

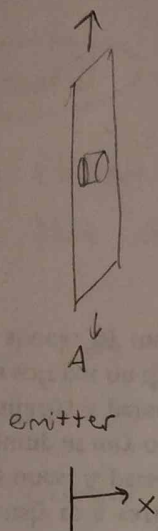
d) The Monopole and dipole parts cancel out, leaving only a quadrupole and higher moments.



d)

$$\vec{F} = \frac{Q^2}{4\pi\epsilon_0} \left[ \frac{-\hat{x}}{\sqrt{(2x_0)^2}} + \frac{-\hat{y}}{\sqrt{(2x_0)^2}} + \frac{\frac{1}{r}(\hat{x}+\hat{y})}{\sqrt{(2x_0)^2 + (2x_0)^2}} \right]$$

3



$\leftarrow d \rightarrow$



$$V(x) = Kx^{4/3}$$

a)  $E = -\vec{\nabla}V = -\hat{x} \frac{\partial}{\partial x} (Kx^{4/3}) = -K \hat{x} \frac{4}{3} x^{1/3}$  [ See part 3.2 for a review ]

b) Consider small box at edge of conductor. We know  $E \perp$  surface of conductor since conductor is an equipotential and  $E \perp$  equipotentials.

Thus  $\int \vec{E} \cdot d\vec{A} \rightarrow \int E dA$ . Also, consider arbitrarily small box, so

$$\int E dA \rightarrow EA$$

- No flux through part of box inside conductor, so have  $EA = \frac{Q_{enc}}{\epsilon_0} = \frac{\sigma A}{\epsilon_0}$

$\Rightarrow \sigma = E \epsilon_0$ , where  $E = E$  infinitesimally close to the conductor;

$$\sigma_e = E(x=0^+) \epsilon_0 \rightarrow = 0$$

\* No discontinuity at surface, so no surface charge.

c)  $\sigma_e = E(x=d^-) \epsilon_0 = \frac{4}{3} K d^{1/3} \epsilon_0$

~~- Here, E would take to be in direction of  $d\vec{A}$ , which is left in this case. To align with x-axis, need to flip sign.~~

$$\Rightarrow \sigma_e = -\frac{4}{3} K d^{1/3} \epsilon_0$$

d)  $\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$   $E = -\vec{\nabla}\phi \Rightarrow \vec{\nabla} \cdot \phi = -\frac{\rho}{\epsilon_0}$

$$\vec{\nabla} \cdot \phi = \frac{\partial}{\partial x} (Kx^{4/3}) = \frac{4}{3} K \frac{\partial}{\partial x} x^{1/3} = \frac{4}{9} K x^{-2/3}$$

$$\Rightarrow \rho(x) = -\frac{4}{9} K \epsilon_0 x^{-2/3}$$

3 cont'd

e) Electrons start with no velocity, and at  $\phi = 0$ . At position  $x$ , they have changed energy by  $q\phi(x) = \frac{1}{2}mv^2$

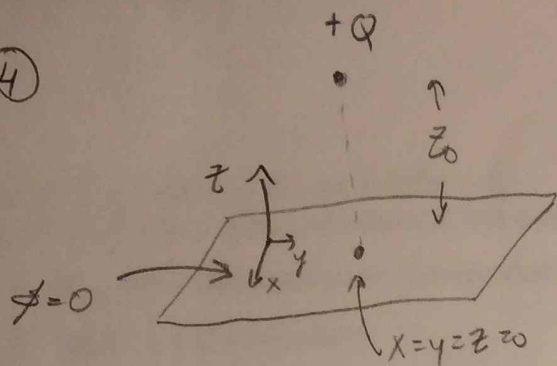
$$\Rightarrow V(x) = \sqrt{\frac{2q}{m}\phi(x)} = \sqrt{\frac{2q}{m}Kx^{4/3}} \quad \text{set.}$$

f) Current density  $\vec{J} \equiv \frac{\Delta(\text{charge})}{\Delta t} \frac{1}{\text{Area}}$

In a cylinder of length  $dx = V(x)dt$ , have  $\rho(x)dx$  charges. These will all pass through the end of the cylinder in time  $dt$ .

$$\begin{aligned} \vec{J}(x) &= \frac{\Delta q}{\Delta t (\text{Area})} \hat{n} = \frac{\rho(x)A dx \hat{n}}{dt A} = \rho(x) \frac{dx}{dt} \hat{n} = \rho(x) \vec{V}(x) \\ &= \left(\frac{-4}{9} K \epsilon_0 x^{-4/3}\right) \sqrt{\frac{2q}{m} K x^{4/3}} \hat{x} = \sqrt{\frac{2q^2}{m} \left(\frac{-4}{9}\right) K \epsilon_0} \hat{x} \end{aligned}$$

4)



Recall Method of images solution:  $\phi(z) = \frac{Q}{4\pi\epsilon_0} \left[ \frac{1}{\sqrt{(z-z_0)^2 + x^2 + y^2}} - \frac{1}{\sqrt{(z+z_0)^2 + x^2 + y^2}} \right]$

a)  $V(z) = Q\phi(z)$  but we only want to include the part of  $\phi(z)$  that isn't from  $Q$  itself: Thus  $V(z) = \frac{-Q}{4\pi\epsilon_0} \frac{1}{\sqrt{(2z)^2}} = \frac{-Q^2}{4\pi\epsilon_0(2z)}$

b)  $T = U_i - U_F = \frac{-Q^2}{8\pi\epsilon_0} \left( \frac{1}{z_0} - \frac{1}{z} \right) = \frac{Q^2}{8\pi\epsilon_0} \left( \frac{1}{z} - \frac{1}{z_0} \right)$

c) At  $z \rightarrow 0$ ,  $T \rightarrow \infty$ . Not physical. This is fine, because at  $z=0$  the method of images analysis fails, because the image charge would have to be at the exact same location as the original charge.

1. (20p) A non-uniform charge distribution.

A non-conducting sphere of radius  $R$  has a non-uniform charge distribution  $\rho(r) = ar$ .

(a) (5p) Find the potential  $\phi(r)$  everywhere.

(b) (5p) Find  $\vec{E}$  everywhere.

(c) (5p) Where is  $\nabla \cdot \vec{E}$  zero and non-zero, and why? (No calculation necessary.)

(d) (5p) If instead the sphere was a conductor with the same amount of total charge  $Q_{tot}$ , how would the electric field change, and why?

2. (20p) Non-conducting shells.

For this problem, we'll look at non-conducting spherical shells and uniformly distributed total charge. Take two such nonconducting shells of uniformly distributed charge, both of radius  $R$ , but with total charges  $Q$ ,  $-Q$ , respectively. The shells are touching. (Remember the shells are non-conducting, so the charges stay put.) See Figure (1).

(a) (5p) Draw the electric field inside the positively charged shell. Can you create the electric field in this shell from a well-placed point charge? If so, how, and if not, why?

(b) (5p) Draw the electric field inside the negatively charged shell. Can you create the electric field in this shell from a well-placed point charge? If so, how, and if not, why?

(c) (5p) Draw the electric field outside both shells. Can you create the electric field outside the shells from well-placed point charges? If so, how, and if not, why?

(d) (5p) What is the work required to separate the shells to infinity? Briefly explain your reasoning.

3. (30p) A simple RC circuit.

The circuit has a switch, a resistor with resistance  $R$ , a battery with electromotive force  $\epsilon$ , and a parallel plate capacitor. At time  $t < 0$ , there is zero charge on the capacitor, and the circuit switch is open. For time  $t \geq 0$ , the switch is closed. See Figure (2).

(a) (5p) Find the charge on the capacitor as a function of time,  $Q(t)$ , and plot  $Q(t)$ .

(b) (5p) Find the current through the circuit as a function of time,  $I(t)$ , and plot  $I(t)$ .

(c) Show that energy is conserved in the circuit.

i. (5p) After a very long time  $t = \infty$ , what was the energy dissipated in the resistor?

ii. (5p) After a very long time  $t = \infty$ , what is the energy stored in the capacitor?

iii. (5p) After a very long time  $t = \infty$ , what is the total work done by the battery?

iv. (5p) Briefly explain energy conservation in the circuit.

4. (30p) Parallel plates and Special Relativity

You have two (approximately infinitely) parallel plates at angle  $\theta = 45$  deg with uniform surface charge density  $\pm\sigma$  (respectively) in the rest frame of the plates (frame  $F$ ). There is an observer in frame  $F'$  moving with velocity  $\vec{v} = -v\hat{x}$  as measured by an observer at rest in frame  $F$ . See Figure(3).

(a) (5p) What is the electric field  $\vec{E}$  (in cartesian coordinates) between the plates the frame  $F$ ?

(b) (5p) What is the electric field  $\vec{E}'$  (in cartesian coordinates) between the plates the frame  $F'$ ?

- (c) (5p) What is the angle of the plates  $\theta'$  in frame  $F'$ ?
- (d) (5p) Draw the electric field from a point charge element of the one of the plates at 3 points  $A, B, C$  in frame  $F$ . See Figure(4).
- (e) (5p) Draw the electric field from a point charge element of the one of the plates at 3 points  $A', B', C'$  in frame  $F'$ . See Figure(4).
- (f) (5p) Is the electric field perpendicular to the plates in frame  $F'$ ? Give a qualitative argument using the results from parts (4b, 4c), and or (4d, 4e).

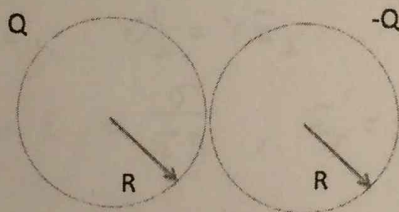
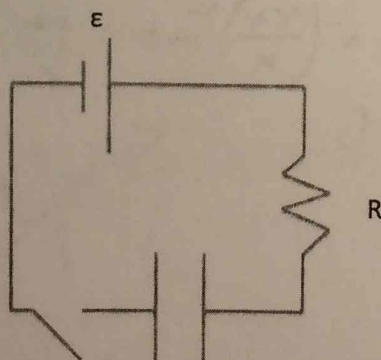


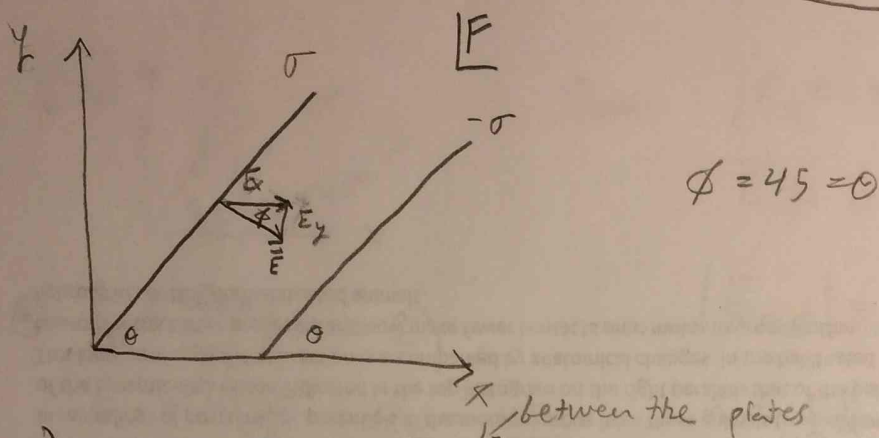
Figure 1. Problem (2).



Switch closed at  $t=0$

Figure 2. Problem (3).

McCurdy Full 2013 Midterm #2, #4



a) Know  $E_{cap} = \frac{\sigma}{\epsilon}$ . Geometry  $\Rightarrow E_x = E \cos \phi = \frac{\sigma}{\sqrt{2} \epsilon_0}$

$\Rightarrow \vec{E} = \frac{\sigma}{\sqrt{2} \epsilon_0} (1, -1)$  between plates. Zero elsewhere

$E_y = -E \sin \phi = \frac{-\sigma}{\sqrt{2} \epsilon_0}$

b) Apply Lorentz transformation for Electric field

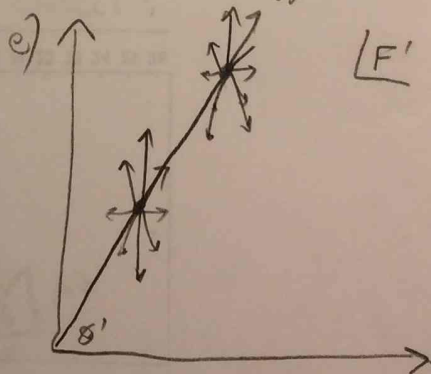
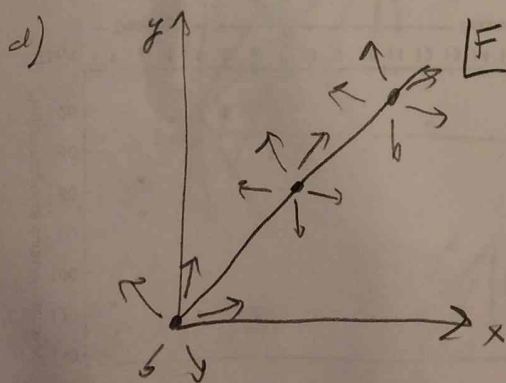
$$E'_{||} = E_{||}, \quad E'_{\perp} = \gamma E_{\perp}$$

$$\Rightarrow E'_x = E_x = \frac{\sigma}{\sqrt{2} \epsilon_0}, \quad E'_y = \gamma E_y = \frac{-\gamma \sigma}{\sqrt{2} \epsilon_0}, \quad \gamma = \frac{1}{\sqrt{1 - (v/c)^2}}$$

c) Apply Lorentz transformation for position (length contraction)

The X-direction gets contracted:  $x' = \frac{x}{\gamma}, \quad y' = y$

$$\Rightarrow \theta' = \tan^{-1}\left(\frac{y'}{x'}\right) = \tan^{-1}\left(\frac{y \gamma}{x}\right) = \tan^{-1}(\gamma \tan \theta) = \tan^{-1}(\gamma)$$



f) No.  $\vec{x}$  and  $\vec{E}$  transform differently