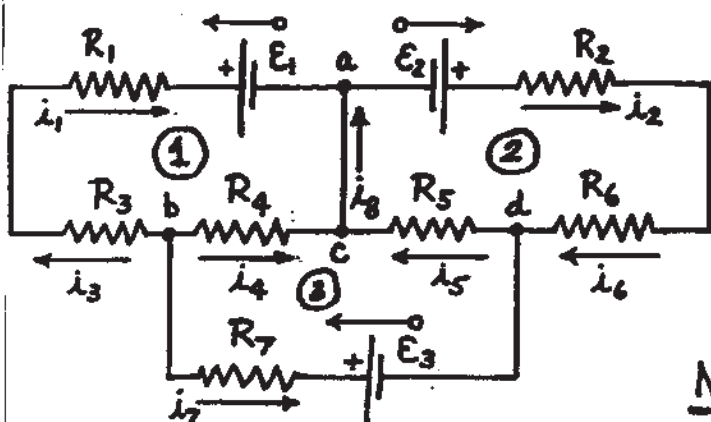


Physics 7B (Sec. 2) Solutions to Midterm #2 April 8, 2003

(1) Number the circuit loops as shown and apply Kirchoff's Loop Theorem, traversing each loop in a clockwise direction, resulting in three equations:



$$(1) -i_1 R_1 - E_1 + i_4 R_4 - i_3 R_3 = 0$$

$$(2) E_2 - i_2 R_2 - i_6 R_6 - i_5 R_5 = 0$$

$$(3) -i_4 R_4 + i_5 R_5 + E_3 + i_7 R_7 = 0$$

Next, apply Kirchoff's Junction Theorem to junctions, a, b, c, d,

obtaining four equations:

$$(4) i_1 + i_8 - i_2 = 0 \quad ; \quad (5) -i_3 - i_7 - i_4 = 0 ;$$

$$(6) i_4 - i_8 + i_5 = 0 \quad ; \quad (7) -i_5 + i_6 + i_7 = 0$$

Equations (1) - (7) are seven equations in eight unknowns, so at least one additional equation is required for solution. A further equation may be obtained from (8)

$$i_3 = i_1 \quad \underline{\text{or}} \quad i_2 = i_6 .$$

Note that i_8 does not appear in the equations for loops ① and ② because i_8 does not pass through a resistor. Thus i_8 does not contribute a potential difference to the Kirchoff loop equations for loops ① and ②.

7B MT #2 SOLUTIONS SPRING 2003 (2) a/d

(2) (a) The capacitance C of parallel plates separated by air - is

$$C = (\epsilon_0 A / d)$$

where the area A of the plates is $A = \pi r^2 = \pi (0.082)^2 = 2.11 \times 10^{-2} \text{ m}^2$ and the separation $d = 1.3 \times 10^{-3}$ meters; $\epsilon_0 = 8.85 \times 10^{-12} \text{ F m}^{-1}$, so

$$C = 1.44 \times 10^{-10} \text{ Farad}$$

The charge Q on either plate is given by $Q = CV$, where the potential difference V between the plates is 120 volts. Then

$$Q = CV = (1.44 \times 10^{-10})(120) = 1.73 \times 10^{-8} \text{ Coulombs} = 17.3 \text{ nC}$$

(b) With a dielectric of dielectric constant K inserted, the capacitance C' is given by

$$C' = KC = (11.7)(1.44 \times 10^{-10}) = 16.8 \times 10^{-10} \text{ Farads}$$

The charge on either plate is still Q because (with battery B removed) there is no conducting path (by which charge might move) between the plates, so

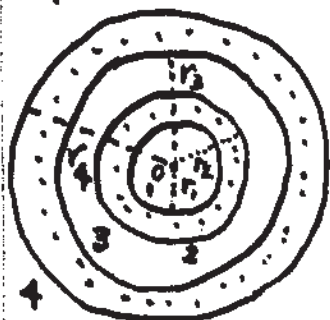
$$Q = C'V' \Rightarrow V' = (Q/C') = \frac{1.73 \times 10^{-8}}{16.8 \times 10^{-10}} = 10.3 \text{ volts}$$

(c) With battery B replaced (now the capacitor has capacitance C') charge can flow (through the battery) between the plates, so

$$Q' = C'V = (16.8 \times 10^{-10})(120) = 2.02 \times 10^{-7} \text{ C.} = 202 \text{ nC}$$

(d) $Q' \neq Q$ because, when the battery B is replaced, charge flows from one plate of the capacitor to the other until (after a long time) the final charge Q' accumulates on either plate.

(3)(a)



Since the shells are conducting, the electric field $|\underline{E}|=0$ inside the shells and all excess charge resides on the surfaces of the shells. An excess charge $(+Q)$ is placed on the inner shell.

To find the charge Q_1 on surface ① (the inner surface of the inner shell), construct a spherical Gaussian surface (concentric with the shells) of radius R , where $r_2 > R > r_1$. This Gaussian surface is inside the inner shell. Since the electric field \underline{E}_1 inside the inner shell is $\underline{E}_1=0$, Gauss' Law becomes

$$\epsilon_0 \int \underline{E}_1 \cdot d\underline{A} = Q_1$$

leading to

$$Q_1 = 0.$$

To find the charge Q_2 on surface ② (the outer surface of the inner shell), note that the problem tells us that there is a charge $(+Q)$ on the inner shell. Since the shell is conducting, all of this charge must be on the outer surface ②, so

$$Q_2 = +Q.$$

To find the charge Q_3 on surface ③ (the inner surface of the outer shell), construct a Gaussian sphere of radius R' , where $r_3 < R' < r_4$, so this Gaussian sphere is inside the outer shell. The electric field \underline{E}_2 inside the outer shell is $\underline{E}_2=0$, so Gauss' Law becomes

$$\epsilon_0 \int \underline{E}_2 \cdot d\underline{A} = Q'$$

(continued \rightarrow)

(3)(a) [continued] where Q' is the total charge enclosed by this Gaussian surface. Therefore

$$Q' = Q_1 + Q_2 + Q_3.$$

Since $\underline{E}_2 = 0$, Gauss' Law gives us $Q' = 0$, so, with $Q_1 = 0$, $Q_2 = +Q$,

$$0 = 0 + Q + Q_3$$

or

$$\boxed{Q_3 = -Q}$$

for the charge on surface (3).

To find the charge Q_4 on surface (4) (the outer surface of the outer shell), we note that the net excess charge Q'' on the outer shell is zero, because excess charge was placed only on the inner shell. Since

$$Q'' = Q_3 + Q_4 = 0$$

and

$$0 = -Q + Q_4$$

$$\boxed{Q_4 = +Q}$$

is the charge on the outer surface (4) of the outer shell.

As a check, note that the total excess charge on the entire system is

$$Q_1 + Q_2 + Q_3 + Q_4 = Q$$

which is the value it should be.

7B MIDTERM #2 SOLUTIONS SPRING 2003 (5) of (11)

(3)[continued] (b) We want to calculate the potential difference $(V_3 - V_2)$ between surfaces ③ and ②, so we need the electric field \underline{E} in the space (vacuum) between surfaces ③ and ②. Since the conductors are spherical and \underline{E} is normal to the conducting surfaces, \underline{E} is radial in direction. We use Gauss' Law to calculate \underline{E} , drawing a spherical Gaussian surface of radius r , where $r_3 > r > r_2$. Gauss' Law is then

$$\epsilon_0 \int \underline{E} \cdot d\mathbf{A} = \epsilon_0 E(r) \int dA = \epsilon_0 E(r) [4\pi r^2] = (Q_1 + Q_2) = Q$$

since $(Q_1 + Q_2) = Q$ is the net charge enclosed by the Gaussian surface. Therefore

$$E(r) = \frac{Q}{4\pi\epsilon_0 r^2}.$$

The potential difference $(V_3 - V_2)$ is given by

$$(V_3 - V_2) = - \int \underline{E} \cdot d\underline{l}$$

taken over $r_3 > r > r_2$. Here $d\underline{l} = (dr)\hat{u}_r$ and the radial $\underline{E}(r) = E(r)\hat{u}_r$, so $(\underline{E} \cdot d\underline{l}) = E(r)dr$ because $(\hat{u}_r \cdot \hat{u}_r) = 1$. Then the integral is

$$(V_3 - V_2) = - \int_{r_2}^{r_3} E(r) dr = - \int_{r_2}^{r_3} \frac{Q}{4\pi\epsilon_0 r^2} = - \frac{Q}{4\pi\epsilon_0} \int_{r_2}^{r_3} \frac{dr}{r^2} = \left. \frac{-Q}{4\pi\epsilon_0} \left| -\frac{1}{r} \right|_{r_2}^{r_3} \right.$$

leading to

$$(V_3 - V_2) = \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{r_3} - \frac{1}{r_2} \right]$$

(c) Since $r_3 > r_2$, $(1/r_3) < (1/r_2)$; since $Q > 0$, $(V_3 - V_2) < 0$ so

$$V_2 > V_3$$

and surface ② is at a higher potential than surface ③. This agrees with $Q_2 = +Q$ and $Q_3 = (-Q)$. (continued \rightarrow)

(3) [continued] (d) To calculate the capacitance C of the system composed of surfaces (2) and (3), we use the definition of capacitance

$$C = (Q/V)$$

which in this case is

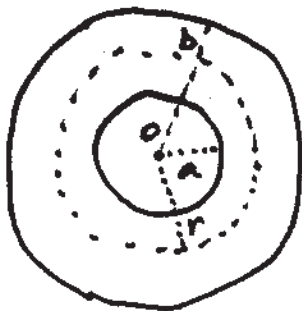
$$C = \frac{Q}{\frac{Q}{4\pi\epsilon_0} \left| \left(\frac{1}{r_3} - \frac{1}{r_2} \right) \right|}$$

$$C = \frac{4\pi\epsilon_0}{\left| \left(\frac{1}{r_3} - \frac{1}{r_2} \right) \right|}$$

as the required capacitance. (We use absolute value of $\left[\frac{1}{r_3} - \frac{1}{r_2} \right]$ because capacitance is always a positive number.)

7B MIDTERM #2 SOLUTIONS SPRING 2003 (7) of (11)

(4)(a) In order to calculate the energy density $u_E(r)$, we need to know the electric field magnitude $|\underline{E}(r)| \equiv E(r)$, which can be calculated using Gauss' Law. Take as the Gaussian surface a sphere of radius r , centered at O . Then Gauss' Law is



$$\epsilon_0 \oint \underline{E} \cdot \hat{n} dA = q_{enc}$$

where q_{enc} is the charge enclosed within the Gaussian surface. To calculate q_{enc} ,

$$q_{enc} = \int \rho(r) dV = \int_a^r \rho_0 r^{-1} [4\pi r^2] dr \quad (\text{where } a \leq r); \rho(r) = \rho_0 r^{-1}$$

where the element of volume dV is a spherical shell of radius r , thickness (dr) and volume $dV = (4\pi r^2) dr$. Then

$$q_{enc} = \int_a^r 4\pi \rho_0 r dr = 4\pi \rho_0 \left[\frac{r^2}{2} \right]_a^r = 4\pi \rho_0 \left(\frac{r^2}{2} - \frac{a^2}{2} \right) = 2\pi \rho_0 (r^2 - a^2)$$

Now that we have q_{enc} , we can use Gauss' Law to calculate $E(r)$. Since $\underline{E}(r)$ is radially outward in direction, $\underline{E}(r) = E(r) \hat{u}_r$. Since $E(r)$ does not depend on direction, we can write, $\hat{n} = \hat{u}_r$,

$$\epsilon_0 \oint E(r) \hat{u}_r \cdot \hat{u}_r dA = \epsilon_0 \oint E(r) dA = \epsilon_0 E(r) \oint dA = q_{enc}$$

so, with $\oint dA = 4\pi r^2$, we obtain

$$\epsilon_0 E(r) [4\pi r^2] = 2\pi \rho_0 (r^2 - a^2)$$

so $E(r) = (\rho_0 / 2\epsilon_0) \left(1 - \frac{a^2}{r^2} \right)$ with which we can calculate the energy density $u_E(r)$

(continued \rightarrow)

(4)(a) [continued] The energy density $u_E(r)$ is then given by

$$u_E(r) = \frac{1}{2} \epsilon_0 |\underline{E}(r)|^2 = \frac{1}{2} \epsilon_0 \left(\rho_0^2 / 4 \epsilon_0^2 \right) \left(1 - \frac{a^2}{r^2} \right)^2$$

$$u_E(r) = \left(\rho_0^2 / 8 \epsilon_0 \right) \left(1 - \frac{a^2}{r^2} \right)^2$$

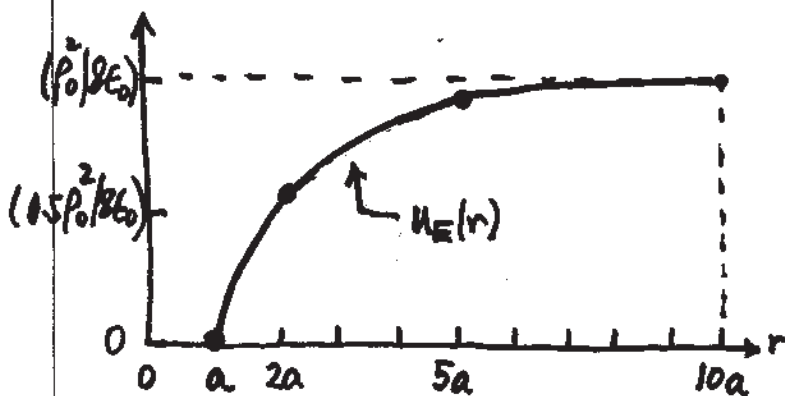
(b) For $b = 10a$, $(b/a) = 10$, $(b^2/a^2) = 100$, $(a^2/b^2) = 0.010$

Then the minimum value $u_E(\text{min})$ occurs for $r = a$: $u_E(\text{min}) = 0$

The maximum value $u_E(\text{max})$ occurs for $r = b$:

$$u_E(\text{max}) = \left(\rho_0^2 / 8 \epsilon_0 \right) \left(1 - \frac{a^2}{b^2} \right)^2 = \left(\rho_0^2 / 8 \epsilon_0 \right) (1.00 - 0.01)^2 \approx \left(\rho_0^2 / 8 \epsilon_0 \right)$$

Plot of $u_E(r)$ as a function of r (for $b = 10a$):



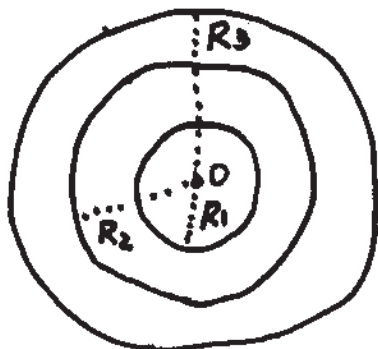
For $r = 2a$: $u_E(2a) = \left(\rho_0^2 / 8 \epsilon_0 \right) \left(1 - \frac{a^2}{4a^2} \right)^2 = \left(\frac{\rho_0^2}{8 \epsilon_0} \right) \left(1 - \frac{1}{4} \right)^2 = (0.56) \frac{\rho_0^2}{8 \epsilon_0}$

For $r = 5a$: $u_E(5a) = \left(\rho_0^2 / 8 \epsilon_0 \right) \left(1 - \frac{a^2}{25a^2} \right)^2 = \left(\frac{\rho_0^2}{8 \epsilon_0} \right) \left(1 - \frac{1}{25} \right)^2 = (0.92) \frac{\rho_0^2}{8 \epsilon_0}$

Units check: $\rho_0^2 = C^2 m^{-4}$, $\epsilon_0 = C^2 N^{-1} m^{-2}$, $\left(\rho_0^2 / \epsilon_0 \right) = N m^{-2} = (N \cdot m) m^{-3} = J m^{-3}$

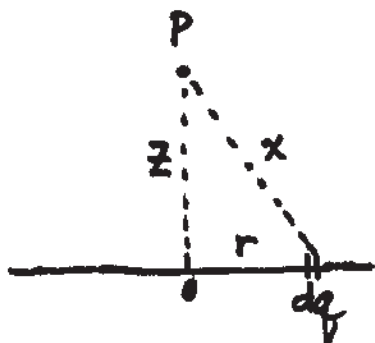
7B MIDTERM #2 SOLUTIONS SPRING 2003 (9) of (11)

(5) Only the disc and the ring matter because the supporting strips are of negligible size. Take as element of area dA an annular ring of radius r and width dr :



TOP VIEW

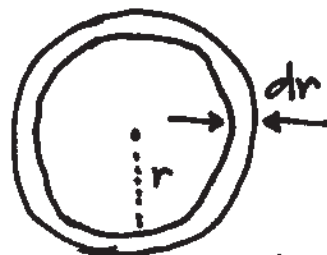
(Point P is a perpendicular distance z above paper.)



SIDE VIEW

(gaps between disc and ring are not shown)

$$x^2 = z^2 + r^2$$



$$dA = d(\pi r^2) = 2\pi r dr$$

Take as the element of charge dq (on ring and disc) the charge on the annular ring:

$$dq = \sigma dA = 2\pi\sigma r dr$$

The element dV of electric potential at point P due to element dq of charge at a distance x from dq is (treating dq as a point charge) given by

$$dV = \frac{1}{4\pi\epsilon_0} \frac{dq}{x} = \frac{2\pi\sigma r dr}{4\pi\epsilon_0 (z^2 + r^2)^{1/2}} \quad \text{where } \begin{cases} 0 \leq r \leq R_1 \\ R_2 \leq r \leq R_3 \end{cases}$$

Since, for a given value of r , all points on an annular ring of radius r are at same distance x from point P, and since, for a given point P, z is fixed, the total electric potential V at point P is given by (since $dV \rightarrow 0$ as $z \rightarrow \infty$) (continued \rightarrow)

(5) [continued]

$$V(z) = \int_{\text{disc}} dV + \int_{\text{ring}} dV$$

so

$$V(z) = \int_0^{R_1} \frac{1}{4\pi\epsilon_0} \frac{2\pi\sigma r dr}{(z^2 + r^2)^{3/2}} + \int_{R_2}^{R_3} \frac{1}{4\pi\epsilon_0} \frac{2\pi\sigma r dr}{(z^2 + r^2)^{3/2}}$$

$$V(z) = \frac{2\pi\sigma}{4\pi\epsilon_0} \left[\int_0^{R_1} \frac{r dr}{(z^2 + r^2)^{3/2}} + \int_{R_2}^{R_3} \frac{r dr}{(z^2 + r^2)^{3/2}} \right]$$

Using one of the integrals supplied,

$$V(z) = \frac{\sigma}{2\epsilon_0} \left\{ \left[(r^2 + z^2)^{-1/2} \right]_0^{R_1} + \left[(r^2 + z^2)^{-1/2} \right]_{R_2}^{R_3} \right\}$$

$$V(z) = \frac{\sigma}{2\epsilon_0} \left[(R_1^2 + z^2)^{-1/2} - z^{-1} + (R_3^2 + z^2)^{-1/2} - (R_2^2 + z^2)^{-1/2} \right]$$

potential at point P; depends on z

(b) The component E_z of the electric field at P in the z -direction (normal to plane of ring and disc) is

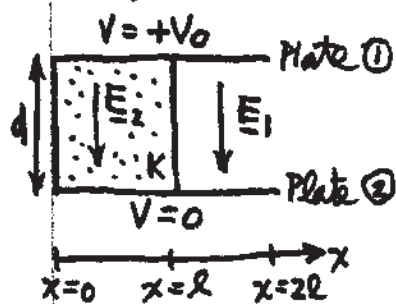
$$E_z = (-dV/dz)$$

so

$$E_z = (-\sigma/2\epsilon_0) \left\{ \frac{z}{(R_1^2 + z^2)^{3/2}} - \frac{1}{z^2} + \frac{z}{(R_3^2 + z^2)^{3/2}} - \frac{z}{(R_2^2 + z^2)^{3/2}} \right\}$$

(c) The negative sign above means that the electric field vector \underline{E}_z at point P points in the direction opposite to the direction in which the potential $V(z)$ increases. Since $V(z)$ decreases with increasing z , so \underline{E}_z points in the direction of increasing z .

(b)(a) The uniform electric field \underline{E}_2 in the dielectric does have the same magnitude as the uniform electric field \underline{E}_1 in the unfilled region of the capacitor. This can be proved as follows. Since the metal plates are equipotential surfaces, the potential difference between the plates is V_0 volts in both regions of the capacitor. From the definition of potential difference,



$$V_0 = -\int \underline{E} \cdot d\underline{l} = -E_2 d = -E_1 d$$

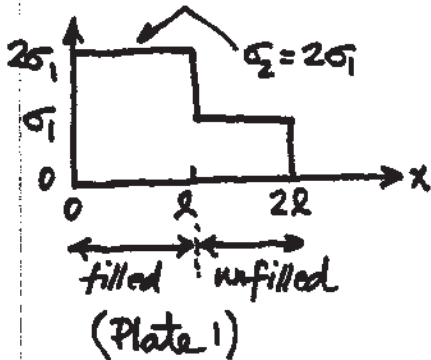
since the potential difference V_0 is the same in both regions of the capacitor. Thus

$$\boxed{E_1 = E_2}$$

(b) We know that the electric field $E_1 = (\sigma_1 / \epsilon_0)$ where σ_1 is the surface charge density on the plates in the unfilled region ($l \leq x \leq 2l$) of the capacitor. Similarly, $E_2 = (\sigma_2 / K\epsilon_0)$ for the filled ($0 \leq x \leq l$) region. Here K is the dielectric constant of the dielectric in the filled region and σ_2 is the surface charge density on the plates in the filled region. Since $E_1 = E_2$, we have

$$(\sigma_1 / \epsilon_0) = (\sigma_2 / K\epsilon_0) \implies \boxed{\sigma_2 = K\sigma_1}$$

so σ_2 is K times σ_1 . For the case $K=2$, $\boxed{\sigma_2 = 2\sigma_1}$ and the required



plot is shown at left. The surface charge density σ_2 on the plates in the filled region is $K (=2)$ times the surface charge density σ_1 on the plates in the unfilled region. Plate 2 has the same surface charge densities of opposite sign.