

(i)(a) The charge $Q(t)$ is given in general by (G.32-14),

$$Q(t) = Q_0 \cos(\omega t + \phi)$$

where Q_0 is the maximum charge on the capacitor. To determine Q_0 , note that $Q(0) = 0$, so

$$0 = Q_0 \cos \phi$$

leading to $0 = \cos \phi \Rightarrow \phi = -\frac{\pi}{2}$

(Note that $Q_0 \neq 0$ because $Q(0) = 0$.) Hence we have

$$Q(t) = Q_0 \cos(\omega t - \frac{\pi}{2})$$

or $Q(t) = Q_0 \sin \omega t$

because $\cos(\omega t - \frac{\pi}{2}) = \sin \omega t$. To determine Q_0 , we are given

$$\left(\frac{dQ}{dt}\right)_{t=0} = I(0) = \omega Q_0 \cos 0 = \omega Q_0 = 2 \text{ Amps}$$

From $L = 3 \times 10^{-3} \text{ H}$, $C = 2.7 \times 10^{-6} \text{ F}$, $\omega = (LC)^{-1/2} = 1.11 \times 10^4 \text{ sec}^{-1}$

$$\Rightarrow \boxed{Q_0 = \frac{2}{\omega} = 1.8 \times 10^{-4} \text{ Coulomb}}$$

(b) We want the maximum value of (dU_E/dt) , where $U_E = (Q^2/2C)$, so

$$U_E(t) = (1/2C) Q_0^2 \sin^2 \omega t \Rightarrow (dU_E/dt) = (Q_0^2/2C) \frac{d}{dt}(\sin^2 \omega t)$$

$$(dU_E/dt) = (Q_0^2/2C)(2\omega) \sin \omega t \cos \omega t = (\omega Q_0^2/2C) \sin(2\omega t)$$

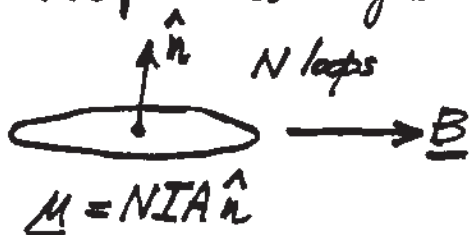
(continued \rightarrow)

(2)(b) [continued] when $X_C = X_L$, $Z = R$, and $\langle P(t) \rangle = \text{max}$.

$$\text{so } \langle P(t) \rangle_{\text{max}} = \frac{\Sigma_0^2 R}{2R^2} = \frac{\Sigma_0^2}{2R} = \frac{(300)^2}{2(5)} = 9000 \text{ Watts}$$

(c) When $X_L = X_C$, $Z = R$, $\cos \phi = (R/Z) = 1$, $\phi = 0$

(3) Coil forms a magnetic dipole of moment $\underline{\mu} = NIA = NIA\hat{n}$



$$\underline{\tau} = \underline{\mu} \times \underline{B}$$

$$\tau = \mu B \sin 90^\circ = NIAB$$

Angle between \hat{n} and $\underline{B} = 90^\circ$

$$\tau = NIAB$$

Since length l of wire is fixed, area A depends on number N of turns in coil. For N loops, circumference of each loop is

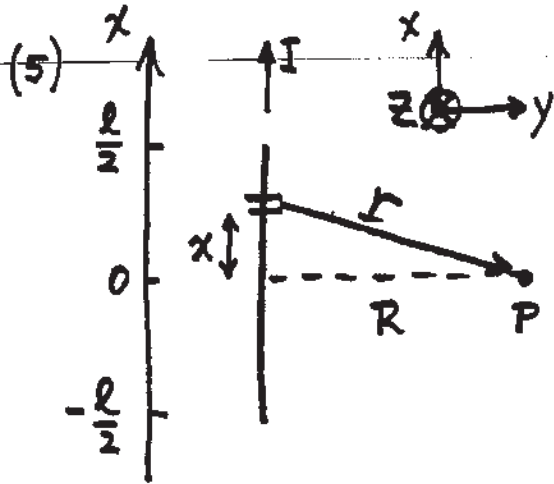
$$c = (l/N) = 2\pi r$$

where $r = (l/2\pi N)$ is radius of each of N loops.

The area A of coil is then $A = \pi r^2 = \pi (l/2\pi N)^2 = \frac{l^2}{4\pi N^2}$

The torque $\tau = NIAB = NIB (l^2/4\pi N^2) = (Il^2 B/4\pi N)$

$$\tau = (Il^2 B/4\pi N)$$



Biot-Savart Law

$$d\vec{B} = \frac{\mu_0 I}{4\pi} \frac{d\vec{l} \times \vec{r}}{r^3}$$

$$x\hat{x} + \vec{r} = R\hat{y}$$

$$\vec{r} = -x\hat{x} + R\hat{y}; \quad r^2 = (x^2 + R^2)$$

$$d\vec{l} = (dx)\hat{x} \quad r^3 = (x^2 + R^2)^{3/2}$$

$$d\vec{B} = \frac{\mu_0 I}{4\pi} \frac{(dx)\hat{x} \times (-x\hat{x} + R\hat{y})}{(x^2 + R^2)^{3/2}} = \frac{\mu_0 I R dx}{4\pi} \frac{(\hat{x} \times \hat{y})}{(x^2 + R^2)^{3/2}}$$

Since $(\hat{x} \times \hat{y}) = \hat{z}$, $d\vec{B} = \left[\frac{\mu_0 I R}{4\pi} \frac{dx}{(x^2 + R^2)^{3/2}} \right] \hat{z}$

Total magnetic field \vec{B} at point P is,

$$\vec{B} = \int_{-l/2}^{l/2} d\vec{B} = \hat{z} \left[\frac{\mu_0 I R}{4\pi} \int_{-l/2}^{l/2} \frac{dx}{(x^2 + R^2)^{3/2}} \right] = \hat{z} \left[\frac{\mu_0 I R}{4\pi} \left. \frac{x}{(x^2 + R^2)^{1/2}} \right|_{-l/2}^{l/2} \right]$$

$$\vec{B} = \hat{z} \left[\frac{\mu_0 I R}{4\pi R^2} \left(\frac{2l}{2\sqrt{\frac{l^2 + 4R^2}{4}}} \right) \right] = \hat{z} \left[\frac{\mu_0 I l}{2\pi R} \frac{1}{\sqrt{l^2 + 4R^2}} \right]$$

$$\boxed{B = \frac{\mu_0 I l}{2\pi R \sqrt{l^2 + 4R^2}}}$$

Physics 57B Solutions to Fourth Exam Aug. 21, 1992

(1)(a) Work done in infinitesimal volume change dV is

$$dW = p dV = \left(\frac{nRT}{V}\right) dV = 4RT \frac{dV}{V}$$

$$\text{Total work } W = \int dW = 4RT \int_{V_i}^{2V_i} \left(\frac{dV}{V}\right) = 4RT \ln 2$$

$$W = 4(8.31)(400)(0.693)$$

$$W = 9216 \text{ Joules}$$

of work done by gas in isothermal expansion.

(b) For ideal gas at constant temperature: $\Delta U = 0 \Rightarrow Q = W$. Since expansion is reversible, $dS = (dQ/T)$ and, since $Q > 0$

$$\Delta S = \int dS = \int \frac{dQ}{T} = \frac{1}{T} \int dQ = \frac{Q}{T} = \frac{9216}{400} = 23 \text{ JK}^{-1}$$

$$\Delta S = 23 \text{ JK}^{-1}$$

(2) Equation of state, with 1 mole, and $T = \text{const.}$,

$$\left(p + \frac{a}{V^2}\right)(V-b) = RT \Rightarrow p(V) = \frac{RT}{(V-b)} - \frac{a}{V^2}$$

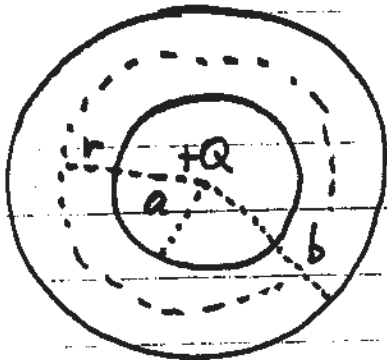
Work W done by gas is

$$W = \int_{V_i}^{V_f} p(V) dV = \int_{V_i}^{V_f} \left[\frac{RT}{(V-b)} - \frac{a}{V^2} \right] dV = RT \int_{V_i}^{V_f} \frac{dV}{(V-b)} - a \int_{V_i}^{V_f} \frac{dV}{V^2}$$

$$W = RT \left[\ln(V-b) \right]_{V_i}^{V_f} + a \left[\frac{1}{V} \right]_{V_i}^{V_f} = RT \ln \left(\frac{V_f - b}{V_i - b} \right) + a \left(\frac{1}{V_f} - \frac{1}{V_i} \right)$$

$$W = RT \ln \left(\frac{V_f - b}{V_i - b} \right) + a \left(\frac{1}{V_f} - \frac{1}{V_i} \right) \text{ work done by gas}$$

(3) Charge density $\rho(r) = Ar^{-1}$ for $a < r < b$; point charge $+Q$ at $r=0$



To calculate $E(r)$, which we can assume spherically symmetric, use Gauss' Law. The Gaussian surface is a sphere of radius r centered at the point charge.

Let $q(r)$ be charge inside Gaussian sphere that is due to the charge density $\rho(r)$, and dq be element of this charge inside a spherical shell of radius r , thickness dr , and volume dV , where $dV = 4\pi r^2 dr$, so

$$dq = \rho(r) [4\pi r^2 dr] \Rightarrow dq = 4\pi A r^{-1} r^2 dr = 4\pi A r dr$$

$$q = \int_a^r dq = 4\pi A \int_a^r r dr = 4\pi A \left[\frac{r^2}{2} \right]_a^r = 2\pi A (r^2 - a^2)$$

Total charge Q' inside Gaussian sphere is $Q' = Q + q(r)$, so

$$Q' = Q + 2\pi A (r^2 - a^2)$$

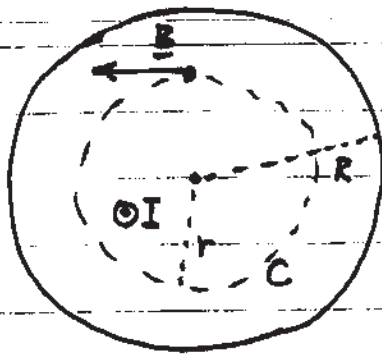
Gauss' Law: $\oint_{G.S.} E(r) dA = \frac{Q'}{\epsilon_0} \Rightarrow \epsilon_0 E(r) \oint dA = Q + 2\pi A (r^2 - a^2)$

$$\epsilon_0 E(r) [4\pi r^2] = [Q + 2\pi A (r^2 - a^2)] \Rightarrow \epsilon_0 E(r) = \frac{Q + 2\pi A (r^2 - a^2)}{4\pi r^2}$$

$$\epsilon_0 E(r) = \frac{Q}{4\pi r^2} + \frac{A}{2} - \frac{2\pi A a^2}{4\pi r^2}$$

$$E(r) = \left(\frac{Q}{4\pi\epsilon_0} - \frac{Aa^2}{2\epsilon_0} \right) \frac{1}{r^2} + \frac{A}{2\epsilon_0} \quad \text{for } (a < r < b)$$

(4)



Let $R =$ radius of circular wire
Current density is

$$J = (I/\pi R^2)$$

Let $i(r) =$ current inside circle of radius r
Since current is spatially uniform,

$$i(r) = J(\pi r^2) = I(r^2/R^2)$$

From Ampere's Law: $\oint_C \underline{B} \cdot d\underline{\ell} = \mu_0 i(r)$ where $i(r)$ is current within C

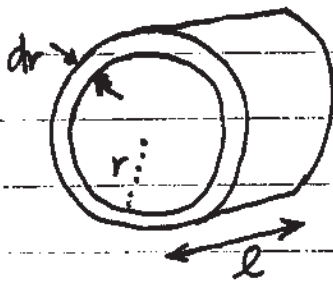
Since $\underline{B} \parallel d\underline{\ell}$, and assuming circular symmetry of $B(r)$, we have

$$\oint_C \underline{B} \cdot d\underline{\ell} = \mu_0 I(r^2/R^2) = B(r) \oint_C d\ell = B(r)[2\pi r] \Rightarrow \boxed{B(r) = (\mu_0 I / 2\pi R^2) r}$$

$$\text{Magnetic energy density } u_B(r) = \frac{1}{2\mu_0} [B(r)]^2 = \frac{1}{2\mu_0} \left(\frac{\mu_0 I^2}{4\pi^2 R^4} \right) r^2$$

To find total magnetic energy U_B within length l of wire

$dU_B = u_B dV$ where $dV = 2\pi r l dr$ is volume of cylindrical shell of length l , radius r , thickness dr



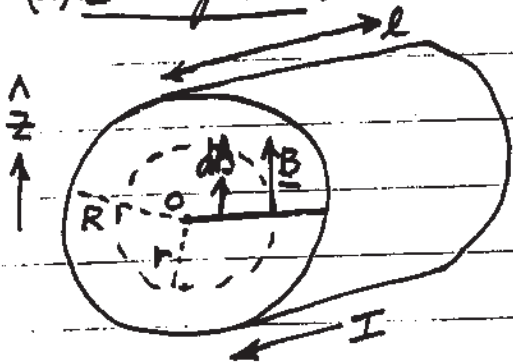
$$dU_B = \frac{1}{2\mu_0} \left(\frac{\mu_0^2 I^2}{4\pi^2 R^4} \right) r^2 (2\pi r l) dr = \frac{\mu_0 I^2 l}{4\pi R^4} r^3 dr$$

$$U_B = \int_0^R dU_B = \frac{\mu_0 I^2 l}{4\pi R^4} \int_0^R r^3 dr = \frac{\mu_0 I^2 l}{4\pi R^4} \left[\frac{R^4}{4} \right]$$

$\frac{U_B}{l}$ is magnetic energy per unit length;

$$\boxed{\frac{U_B}{l} = \frac{\mu_0 I^2}{16\pi}}$$

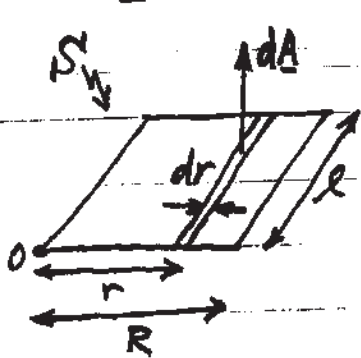
(3) Diagram is below: $\underline{B}(r) = \frac{\mu_0 i(r) \hat{z}}{2\pi r}$



where $i(r)$ is current inside cylinder of radius r and length l ; $\underline{B}(r)$ is field (due to $i(r)$) at distance r from center O of wire. Since current I is spatially uniform, $i(r) = I(r^2/R^2)$,

$$\text{so } \underline{B}(r) = \frac{\mu_0 I r^2}{2\pi R^2 r} = \left(\frac{\mu_0 I}{2\pi R^2}\right) r$$

Let $d\underline{A}$ be vector element of area in plane S' in drawing given, where



$$\left. \begin{aligned} d\underline{A} &= (dA)\hat{z} = (l dr)\hat{z} \\ \underline{B}(r) &= \left(\frac{\mu_0 I}{2\pi R^2} r\right)\hat{z} \end{aligned} \right\} \underline{B} \cdot d\underline{A} = \frac{\mu_0 I l}{2\pi R^2} r dr$$

Total flux Φ through surface S' is $\Phi = \int d\Phi$

$$\Phi = \int d\Phi = \int_0^R \underline{B} \cdot d\underline{A} = \frac{\mu_0 I l}{2\pi R^2} \int_0^R r dr = \frac{\mu_0 I l}{2\pi R^2} \left[\frac{R^2}{2}\right]$$

$$\boxed{\Phi = \frac{\mu_0 I l}{4\pi}} \text{ is total flux through } S'$$