

Problem 1

Let T be the tension in the rope between the drum and the pulley, R be the radius of both the drum and the pulley, and $M = 2\text{ kg}$ be the mass of both the drum and the pulley.

$$\Sigma\tau_{drum} = TR = I_{drum}\alpha_{drum} = MR^2\alpha_{drum}$$

$$\Sigma\tau_{pulley} = TR - FR = I_{pulley}\alpha_{pulley} = \frac{1}{2}MR^2\alpha_{pulley}$$

$$F_{pulley} = T + F - Mg = ma_{pulley} = 0$$

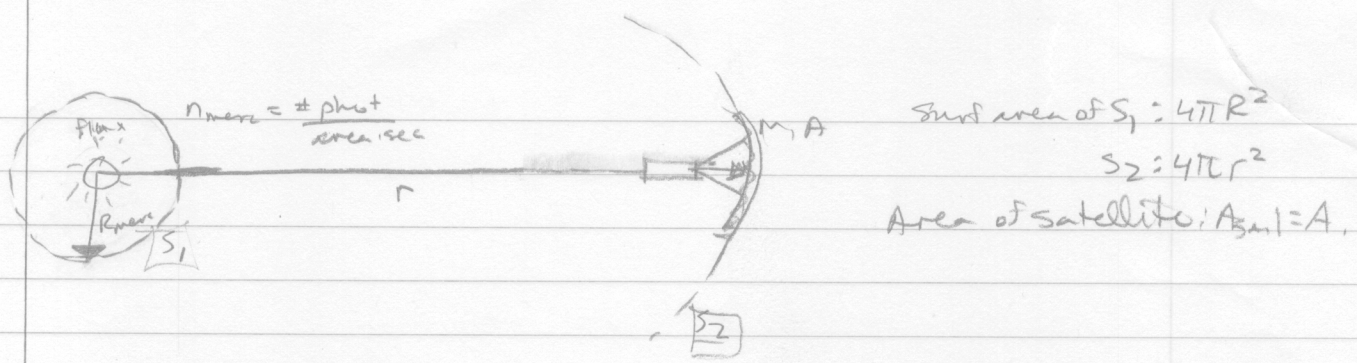
$$\alpha_{drum} = -\alpha_{pulley} \Rightarrow TR = MR^2\alpha \text{ and } FR - TR = \frac{1}{2}MR^2\alpha$$

$$TR = 2FR - 2TR \Rightarrow 3T = 2F \Rightarrow T = \frac{2}{3}F$$

$$T + F - mg = 0 \Rightarrow \frac{2}{3}F + F = \frac{5}{3}F = mg \Rightarrow F = \frac{3}{5}mg$$

$$F = \frac{3}{5}(2\text{ kg})(9.8 \frac{\text{m}}{\text{s}^2}) = 11.76 \text{ N}$$

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a) $F_g = \frac{GMm}{r^2}$ ← $F_{\text{change dir. of photons}} \approx \frac{\Delta p_{\text{phot}}}{\Delta t}$ ← perfect reflector; $\Delta p = 2p$ change in mom. spacecraft.

photons $\frac{= n_s}{\text{area} \cdot \text{sec}} \Rightarrow \text{photons/sec} = n_{merc} \cdot A_{merc} = n_{merc} \cdot 4\pi R^2$
sphere (S_1)

At mercury: $\frac{\text{Total momentum}}{\text{sec}} = \frac{\# \text{ photons}}{\text{sec}} \cdot \frac{\text{mom}}{\text{photon}} = 4\pi R_{merc}^2 n_{merc} p$
= total mom/sec through sphere at radius r.

At the large sphere:
 $\frac{\text{mom}}{\text{area} \cdot \text{sec}} = \frac{(4\pi R_{merc}^2 n_{merc} p)}{4\pi r^2} = \frac{R^2}{r^2} n_{merc} p$

⇒ Momentum hitting the sail per second: $\frac{R^2}{r^2} n_{merc} p \cdot A_{sail}$

⇒ $\frac{\Delta p}{\Delta t}$ (change in mom. per sec):
 $= 2 \left(\frac{R^2}{r^2} n_{merc} p \cdot A \right) = F_{\text{photon}}$
perfect reflector (see above)

$\Sigma F = F_{\text{grav}} + F_{\text{photon}} = -\frac{GM_S m}{r^2} + \frac{2R^2}{r^2} n_{merc} p \cdot A$

⇒ $\Sigma \vec{F} = \frac{2R^2 n_{merc} p \cdot A}{r^2} - GM_S m \hat{r}$

b) Max attainable speed? Begin $r = R_{merc}$.

$\frac{1}{2} m v_{\text{max}}^2 - \frac{1}{2} m v_0^2 = \int_{R_{merc}}^{\infty} \vec{F} \cdot d\vec{r} = \int_{R_{merc}}^{\infty} (2p A n_{merc} R_{merc}^2 - GM_S m) \frac{1}{r^2} dr$

⇒ $\frac{1}{2} m v_{\text{max}}^2 - \frac{1}{2} m v_0^2 = \left(\frac{2p A n_{merc} R_{merc}^2}{r} - \frac{GM_S m}{r} \right) \Big|_{R_{merc}}^{\infty}$

Recall $\frac{v_0^2}{r} = \frac{GM_{sun}}{R_{merc}}$ (assuming circular orbit for mercury, to match mize/get expression for v_0^2)

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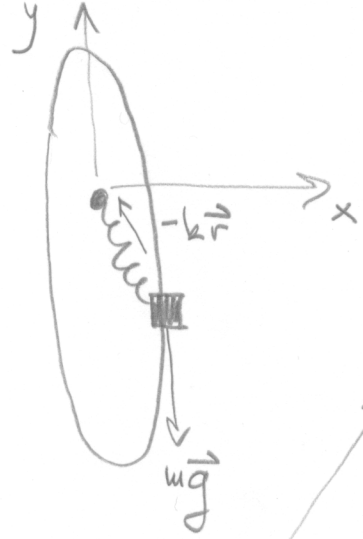
$$\text{Then } v_0^2 = \frac{GM_{\text{sun}}}{R_{\text{merc}}} \quad \text{at } r = R_{\text{merc}}$$

$$\Rightarrow v_{\text{max}}^2 = \frac{2(2\rho A R_{\text{merc}} R_{\text{merc}}^2 - GM_{\text{sun}})}{m R_{\text{merc}}} + \frac{2m v_0^2}{2m}$$

$$= \frac{4\rho A R_{\text{merc}} R_{\text{merc}}}{m} - \frac{2GM_{\text{sun}}}{R_{\text{merc}}} + \frac{GM_{\text{sun}}}{R_{\text{merc}}}$$

$$v_{\text{max}} = \sqrt{\frac{4\rho A R_{\text{merc}} R_{\text{merc}}}{m} - \frac{GM_{\text{sun}}}{R_{\text{merc}}}} \quad \text{1/1}$$

3



$$m\vec{a} = \vec{F}_{net} = -k\vec{r} + m\vec{g}$$

$$ma_x = m \frac{d^2x}{dt^2} = -kx$$

$$may = m \frac{d^2y}{dt^2} = -ky - mg$$

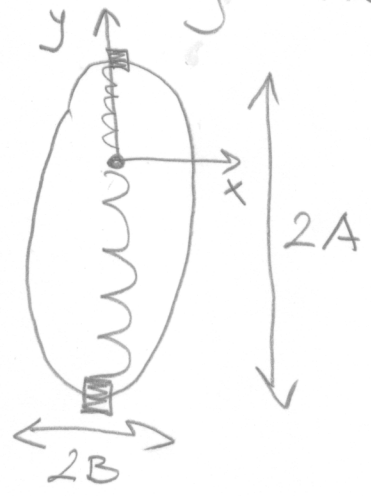
$$y(t) = A \cos(\omega t) - \frac{mg}{k}$$

$$\frac{dy}{dt}(t) = -\omega A \sin(\omega t)$$

$$\frac{d^2y}{dt^2}(t) = -\omega^2 A \cos(\omega t)$$

satisfies equation
with $\omega = \sqrt{\frac{k}{m}}$

Similarly $x(t) = -B \sin(\omega t)$



but we don't need this

$$R_A = \max y(t) \quad (\text{when } \cos = 1)$$

$$R_A = A - \frac{mg}{k}$$

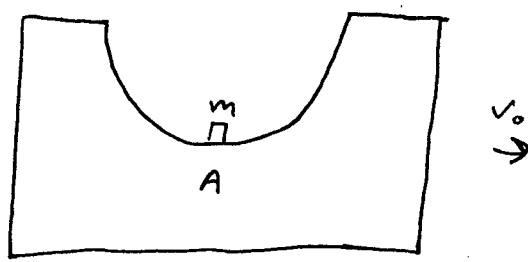
$$R_B = |\min y(t)| \quad (\text{when } \cos = -1)$$

$$R_B = A + \frac{mg}{k}$$

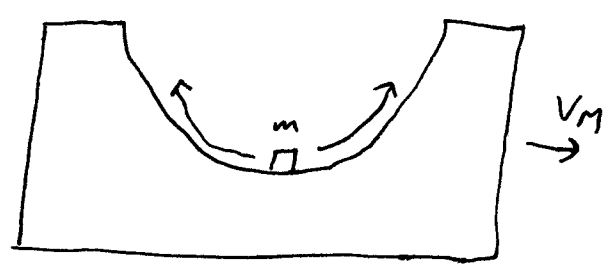
We don't know
A or B but we
can solve
this problem:

$$R_B - R_A = A + \frac{mg}{k} - \left(A - \frac{mg}{k} \right) = \frac{2mg}{k}$$

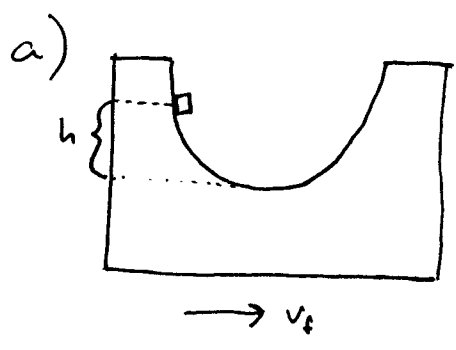
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$t=0$



$t>0$



Momentum is conserved, and when the small block reaches its maximum height it is at rest w.r.t. the big block, so both blocks are moving with velocity v_f .

$$p_0 = p_f$$

$$Mv_0 = (M+m)v_f$$

$$v_f = \frac{M}{M+m} v_0$$

Energy is conserved:

$$\frac{1}{2} Mv_0^2 = \frac{1}{2} (M+m)v_f^2 + mgh = \frac{1}{2} \frac{M^2 v_0^2}{M+m} + mgh$$

$$h = \frac{v_0^2 M}{2mg} \left(1 - \frac{M}{M+m}\right) = \boxed{\frac{v_0^2}{2g} \left(\frac{M}{M+m}\right)}$$

b) In CM frame, the small mass moves to the left and rises to h , then falls back down and reaches its maximum velocity at A.

Momentum is conserved

$$Mv_0 = mv_m + Mv_M$$

$$v_M = \frac{Mv_0 - mv_m}{M}$$

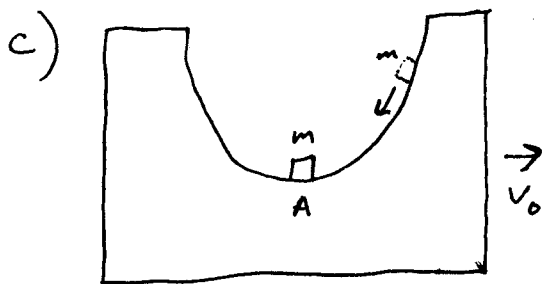
Energy is conserved

$$\frac{1}{2} M v_0^2 = \frac{1}{2} m v_m^2 + \frac{1}{2} M v_M^2 = \frac{1}{2} m v_m^2 + \frac{1}{2} M \left(\frac{M v_0 - m v_m}{M} \right)^2$$

$$M v_0^2 = m v_m^2 + M v_0^2 + \frac{m^2}{M} v_m^2 - 2 m v_0 v_m$$

$$v_m \left(1 + \frac{m}{M} \right) = 2 v_0$$

$$v_m = \frac{2 v_0}{1 + \frac{m}{M}}$$



When m approaches A from the right, the system will be in its initial state. This happens for $\boxed{\text{odd } N}$

Problem 5

Conservation of momentum

$$m_a \vec{v}_0 = m_a \vec{v}_{fa} + m_b \vec{v}_{fb}$$

This is an elastic collision, so we can conserve kinetic energy

$$\frac{1}{2} m_a v_0^2 = \frac{1}{2} m_a v_{fa}^2 + \frac{1}{2} m_b v_{fb}^2$$

For $m_a = m_b$

$$\begin{aligned}\vec{v}_0 &= \vec{v}_{fa} + \vec{v}_{fb} \text{ and } v_0^2 = v_{fa}^2 + v_{fb}^2 \\ \vec{v}_0 \cdot \vec{v}_0 &= (\vec{v}_{fa} + \vec{v}_{fb}) \cdot (\vec{v}_{fa} + \vec{v}_{fb}) \\ \Rightarrow v_0^2 &= v_{fa}^2 + v_{fb}^2 + 2\vec{v}_{fa} \cdot \vec{v}_{fb} \\ \Rightarrow v_{fa}^2 + v_{fb}^2 &= v_{fa}^2 + v_{fb}^2 + 2\vec{v}_{fa} \cdot \vec{v}_{fb} \\ \Rightarrow \vec{v}_{fa} \cdot \vec{v}_{fb} &= 0\end{aligned}$$

Since $\vec{v}_{fa} \cdot \vec{v}_{fb} = 0$, the angle between \vec{v}_{fa} and \vec{v}_{fb} is 90° ■